

Aspects of Quantum Cosmology

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November 6, 2013

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- Amalgamation of Quantum Theory and Cosmology.

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- The very early universe before the Planck Time scale when the length of the universe is order of Planck length.

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- The very early universe before the Planck Time scale when the length of the universe is order of Planck length.
- The thermodynamical arrow of time and its matching with cosmological arrow of time.

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- We need to find a parameter which is timelike against which our universe will evolve.
- We need to build up a configuration space for universe where we will define Quantum Dynamics.
- **The definition of Probability is not so well formulated. Some anisotropic models reveal time dependent norm when we define probability in standard way**
- **Our aim is to look for a suitable definition of probability so as to alleviate this time dependence leading to Unitarity of evolution operator**

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Define extrinsic curvature of the hypersurface to be

$$K_{ij} = -n_{i;j} = \frac{1}{2N} \left(N_{i|j} + N_{j|i} - \frac{\partial h_{ij}}{\partial t} \right) \quad (2)$$

where n_i is normal to the spacelike hypersurface.

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where n_i is normal to the spacelike hypersurface.

$$S = \frac{1}{4\kappa^2} \left[\int_{\mathcal{M}} d^4x \sqrt{-g} R_{(4)} + 2 \int_{\partial\mathcal{M}} d^3x \sqrt{h} K \right] + S_{matter} \quad (3)$$

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$$\pi^0 \equiv \frac{\partial L}{\partial \dot{N}} = 0 \quad (6)$$

$$\pi^i \equiv \frac{\partial L}{\partial \dot{N}_i} = 0 \quad (7)$$

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where

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Introducing DeWitt Metric

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- If we consider a phase space described by h_{ij} and π^{ij} ; \mathcal{G}_{ijkl} will provide a metric on that space

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$$Riem(\Sigma) \equiv \{h_{ij}, \Phi(x) : x \in \Sigma\} \quad (15)$$

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Hence **Superspace** is defined to be

$$Superspace(\Sigma) \equiv \frac{Riem(\Sigma)}{Diff(\Sigma)} \quad (16)$$

Metrizing the Superspace

The DeWitt metric provides a metric on Superspace, given by

$$\mathcal{G}_{AB}(x) = \mathcal{G}_{(ij)(kl)}(x) \quad (17)$$

where $A, B \in \{h_{ij} : i \leq j, i, j \in \{1, 2, 3\}\}$

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- The inverse metric is given by \mathcal{G}^{AB} so that

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Superspace for Homogeneous Closed Universe

$$S = \int dt \left[\frac{1}{2N} \mathcal{G}_{AB} \dot{q}^A \dot{q}^B - N U(q) \right] \quad (20)$$

where $U(q) = V(\Phi) - \frac{R_{(3)}}{4\kappa^2}$

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$$H = N \left[\frac{1}{2} \mathcal{G}^{AB} \pi_A \pi_B + U(q) \right] \quad (21)$$

$$\left[\frac{1}{2} \mathcal{G}^{AB} \pi_A \pi_B + U(q) \right] = 0 \quad (22)$$

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The conserved current is

$$j^\alpha = \frac{-i\hbar}{2} \mathcal{G}^{\alpha\beta} (\psi^* \nabla_\beta \psi - \psi \nabla_\beta \psi^*) \quad (25)$$

Probability

We take ansatz $\psi = A(q)e^{i\frac{S(q)}{\hbar}}$ to arrive

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Hence probability is defined to be

$$P = \int j^\alpha d\Sigma_\alpha \quad (27)$$

Small Quantum Subsystem

Consider a small quantum subsystem with $H_{q'}$ being the hamiltonian:

$$ds^2 = \mathcal{G}_{\alpha\beta} q^\alpha q^\beta + \mathcal{G}_{\alpha m} q^\alpha q'^m + \mathcal{G}_{mn} q'^m q'^n \quad (28)$$

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$$\mathcal{G}_{\alpha m}(q, q') = O(\lambda) \quad (30)$$

Arriving the Schrodinger Equation for Subsystem

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$$NH_{q'}\chi = i\hbar \frac{\partial \chi}{\partial t} \quad (32)$$

Conserved Currents

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Normalisation condition yields

$$\int j_0^\alpha d\Sigma_\alpha \int \rho_\chi d\Omega_{q'} = 1 \quad (33)$$

- This sets normalisation condition on the probability of small quantum subsystem.

Probability conservation in Subsystem

$$\nabla_\alpha j^\alpha + \nabla_m j^m = 0 \quad (34)$$

- Write $j^\alpha = (|A|^2 \hbar \nabla^\alpha S) \rho_\chi$ and use $\nabla_\alpha j_0^\alpha = 0$

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- Everything falls into place and we get back our standard QM result!!

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- The anisotropic Models, Bianchi-V and Bianchi-IX show nonunitarity upon a particular choice of time parameter.
- The DeWitt metric always has a signature of $(-, +, +..)$. So we will try to take coordinate with opposite signature as time.
- In that case the resulting equation may not be first order in time derivative unlike Schrodinger equation.

References

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Acknowledgement

The speaker acknowledges a debt of gratitude to

- Prof. N. Banerjee, Dept. of Physical Sciences, IISER Kolkata for his valuable time and guidance.

THANK YOU